

Chapter 34

(continued)

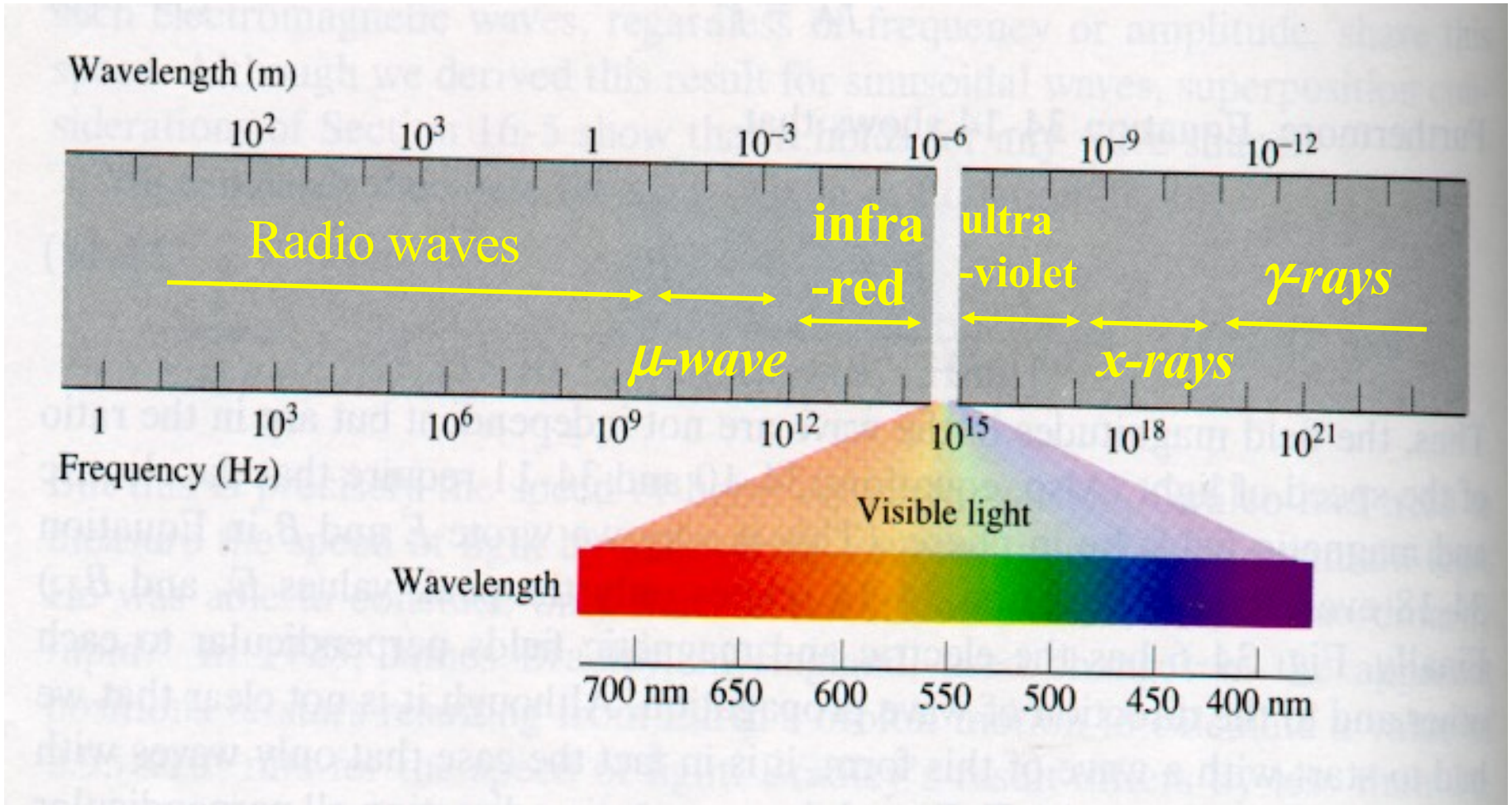
The Laws of Electromagnetism

Maxwell's Equations

Displacement Current

Electromagnetic Radiation

The Electromagnetic Spectrum



Maxwell's Equations of Electromagnetism in Vacuum (no charges, no masses)

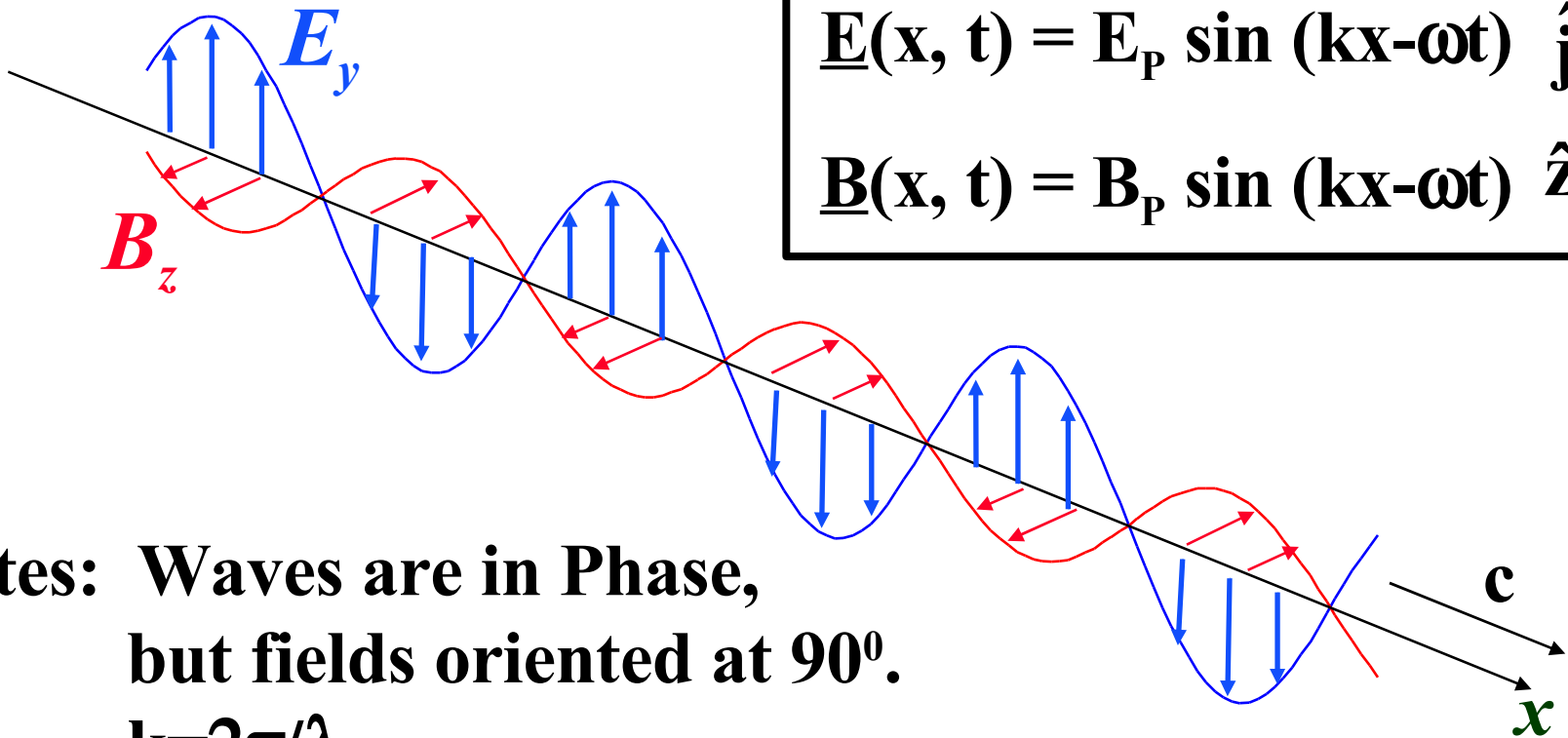
$$\oint \underline{E} \cdot \underline{dA} = 0$$

$$\oint \underline{B} \cdot \underline{dA} = 0$$

$$\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Plane Electromagnetic Waves



**Notes: Waves are in Phase,
but fields oriented at 90° .**

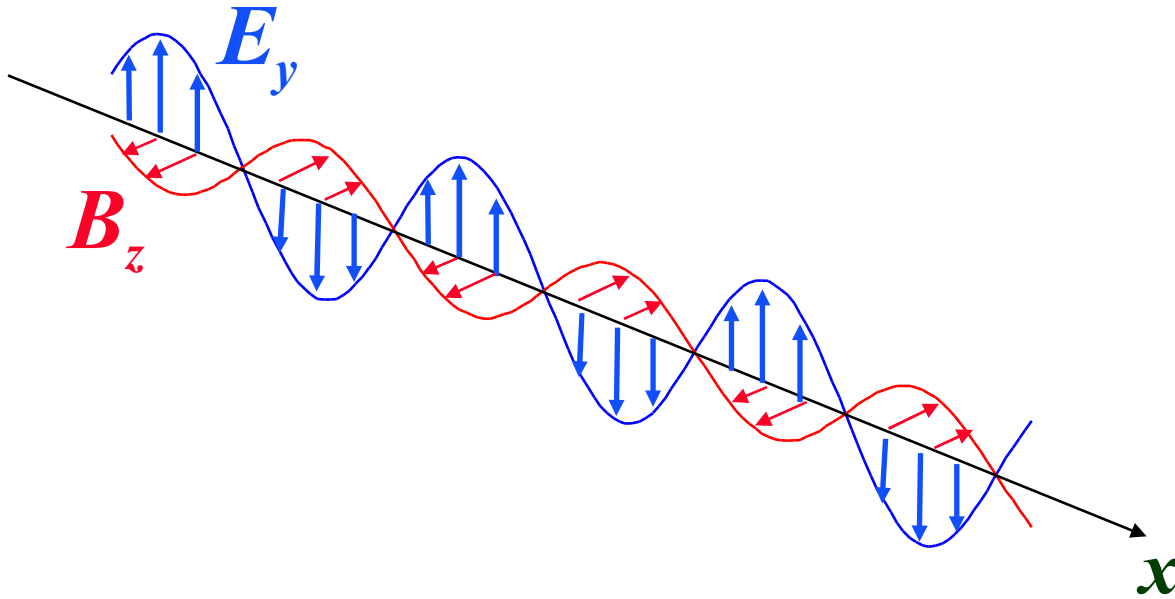
$$k = 2\pi/\lambda.$$

Speed of wave is $c = \omega/k$ ($= f\lambda$)

$$c = 1 / \sqrt{\epsilon_0 \mu_0} = 3 \times 10^8 \text{ m/s}$$

At all times $E = cB$.

Plane Electromagnetic Waves



Note: $\sin(\omega t - kx) = -\sin(kx - \omega t) \Rightarrow$ notations are interchangeable. $\sin(\omega t - kx)$ and $\sin(kx - \omega t)$ represent waves traveling towards +x, while $\sin(\omega t + kx)$ travels towards -x.

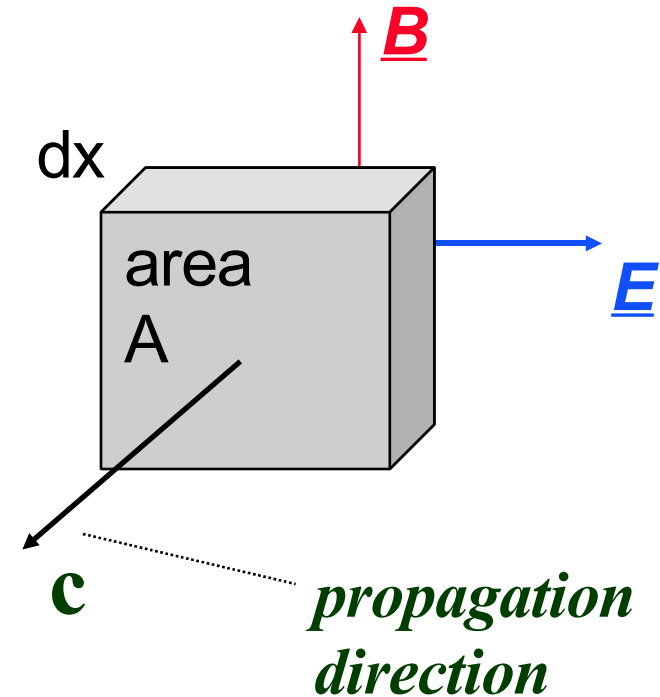
Energy in Electromagnetic Waves

- Electric and magnetic fields contain energy, potential energy stored in the field: u_E and u_B
 u_E : $\frac{1}{2} \epsilon_0 E^2$ electric field energy density
 u_B : $(1/\mu_0) B^2$ magnetic field energy density
- The energy is put into the oscillating fields by the sources that generate them.
- This energy can then propagate to locations far away, at the velocity of light.

Energy in Electromagnetic Waves

Energy per unit volume is

$$u = u_E + u_B = \frac{1}{2}(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2)$$



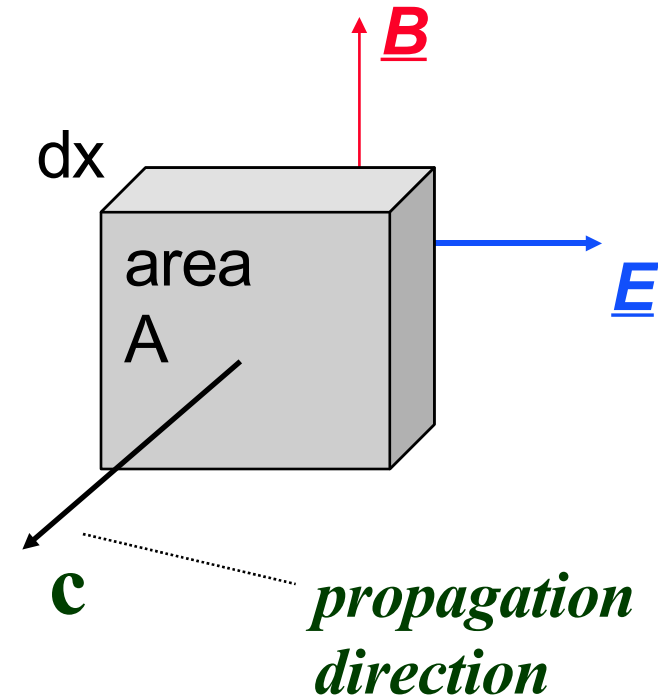
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Thus the energy, dU , in a box of area A and length dx is

$$dU = \frac{1}{2}(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2) A dx$$



Energy in Electromagnetic Waves

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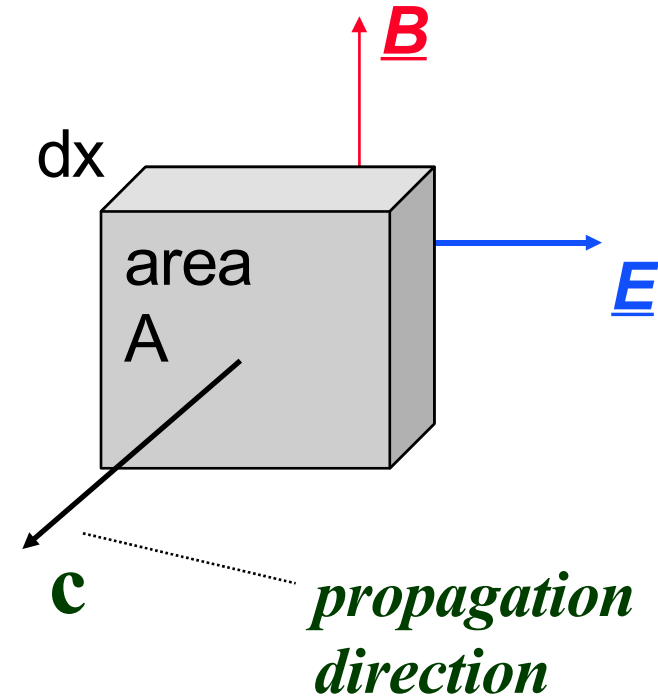
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$$dU = \frac{1}{2}(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2) A dx$$

Let the length dx equal cdt . Then all of this energy leaves the box in time dt . Thus energy flows at the rate

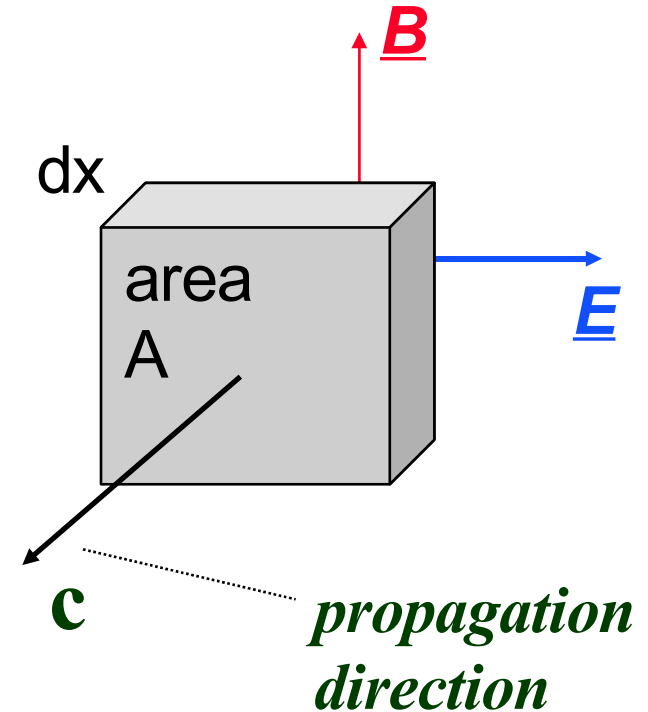
$$\frac{dU}{dt} = \frac{1}{2}(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2) A c$$



Energy Flow in Electromagnetic Waves

Rate of energy flow:

$$\frac{dU}{dt} = \frac{c}{2} (\epsilon_0 E^2 + \frac{1}{\mu_0} B^2) A$$



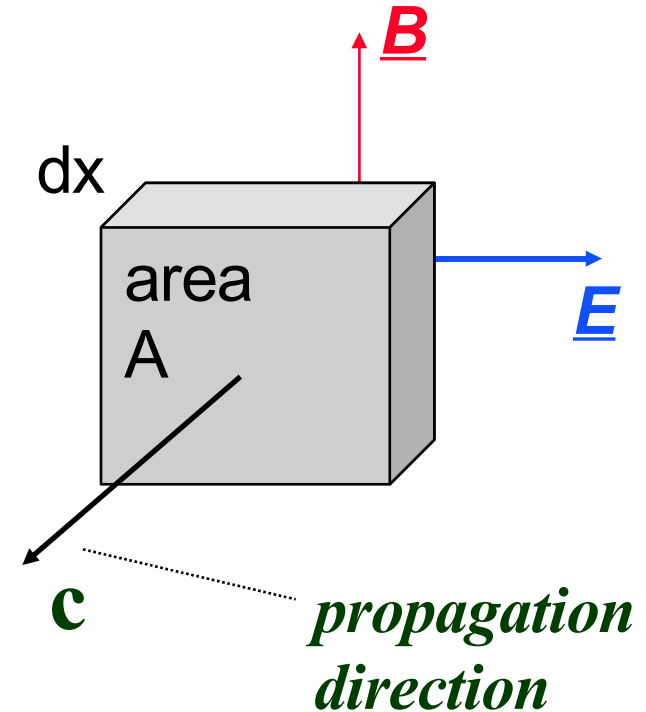
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*We define the **intensity S** , as the rate of energy flow per unit area:*

$$S = \frac{c}{2} (\epsilon_0 E^2 + \frac{1}{\mu_0} B^2)$$



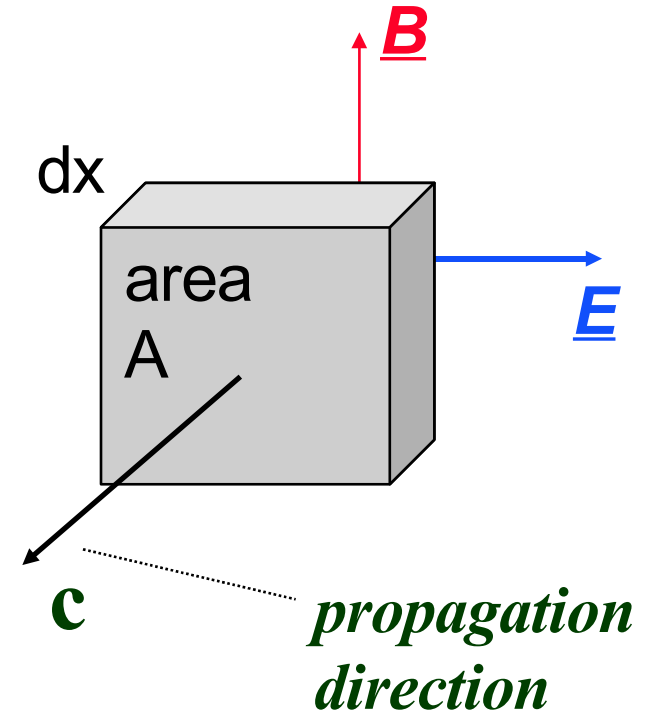
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Rearranging by substituting $E=cB$ and $B=E/c$, we get,

$$S = \frac{c}{2} (\epsilon_0 c E B + \frac{1}{\mu_0 c} E B) = \frac{1}{2\mu_0} (\epsilon_0 \mu_0 c^2 + 1) E B = \frac{E B}{\mu_0}$$

The Poynting Vector

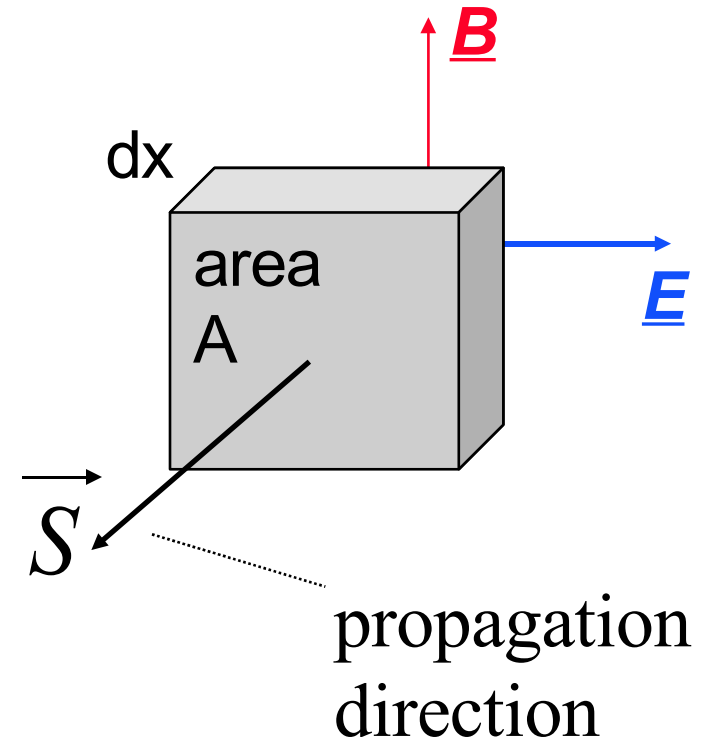
In general, we find:

$$\underline{S} = (1/\mu_0) \underline{E} \times \underline{B}$$

S is a vector that points in the direction of propagation of the wave and represents the rate of energy flow per unit area.

We call this the “**Poynting vector**”.

Units of S are $Jm^{-2} s^{-1}$, or Watts/m².

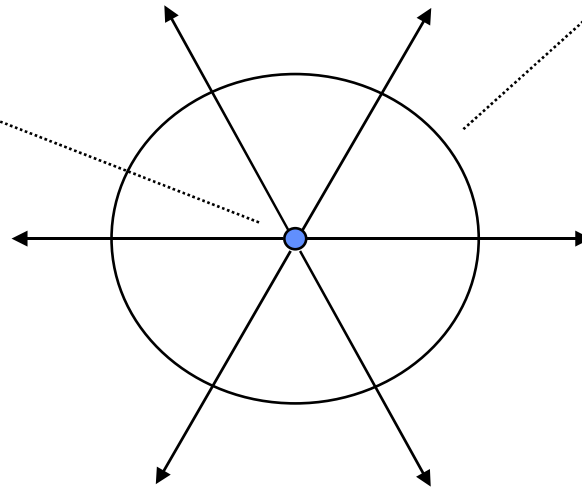


The Inverse-Square Dependence of \bar{S}

A point source of light, or any radiation, spreads out in all directions:

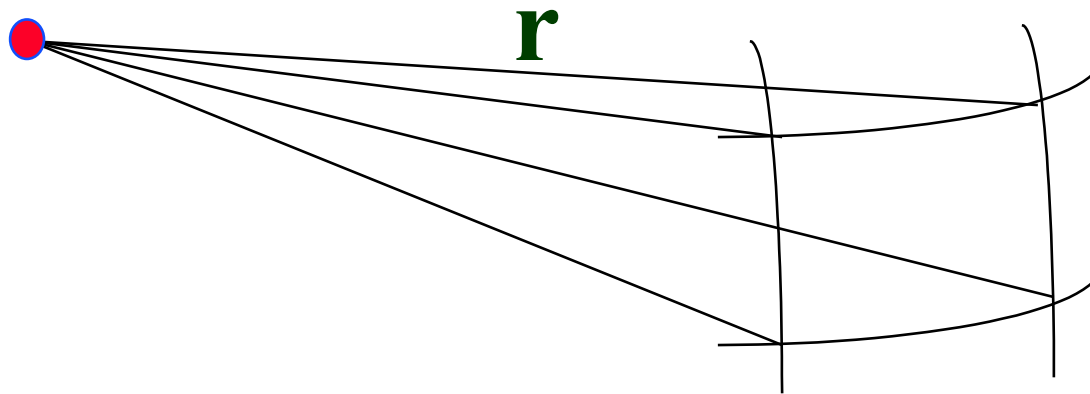
Source

$$\bar{S} = \frac{\bar{P}}{4\pi r^2}$$



Power, P , flowing through sphere is same for any radius.

Source



$$\text{Area} \propto r^2$$

$$\bar{S} \propto \frac{1}{r^2}$$

Example:

An observer is 1.8 m from a point light source whose average power $P = 250 \text{ W}$. Calculate the rms fields in the position of the observer.

Intensity of light at a distance r is $S = P / 4\pi r^2$

$$I = \frac{P}{4\pi r^2} = \frac{1}{\mu_0 c} E_{rms}^2$$

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$$\therefore E_{rms} = \sqrt{\frac{P\mu_0 c}{4\pi r^2}} = \sqrt{\frac{(250\text{W})(4\pi 10^{-7} \text{H/m})(3 \cdot 10^8 \text{m/s})}{4\pi (1.8\text{m})^2}}$$

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$$\therefore E_{rms} = 48\text{V/m}$$

$$\therefore B = \frac{E_{rms}}{c} = \frac{48\text{V/m}}{3.10^8 \text{m/s}} = 0.16\mu\text{T}$$

Wave Momentum and Radiation Pressure

Momentum and energy of a wave are related by, $p = U / c$.

Now, Force = $d p / dt = (dU/dt)/c$

pressure (radiation) = Force / unit area

$$P = (dU/dt) / (A c) = S / c$$

$$\therefore \text{Radiation Pressure} \Rightarrow P_{rad} = \frac{\bar{S}}{c}$$

Example: Serious proposals have been made to “sail” spacecraft to the outer solar system using the pressure of sunlight. How much sail area must a 1000 kg spacecraft have if its acceleration is to be 1 m/s^2 at the Earth’s orbit? Make the sail reflective.

Can ignore gravity.

Need $F=ma=(1000\text{kg})(1 \text{ m/s}^2)=1000 \text{ N}$

This comes from pressure: $F=PA$, so $A=F/P$.

Here P is the radiation pressure of sunlight:

Sun’s power = $4 \times 10^{26} \text{ W}$, so $S=\text{power}/(4\pi r^2)$ gives

$$S = (4 \times 10^{26} \text{ W}) / (4\pi(1.5 \times 10^{11} \text{ m})^2) = 1.4 \text{ kW/m}^2.$$

Thus the pressure due to this light, reflected, is:

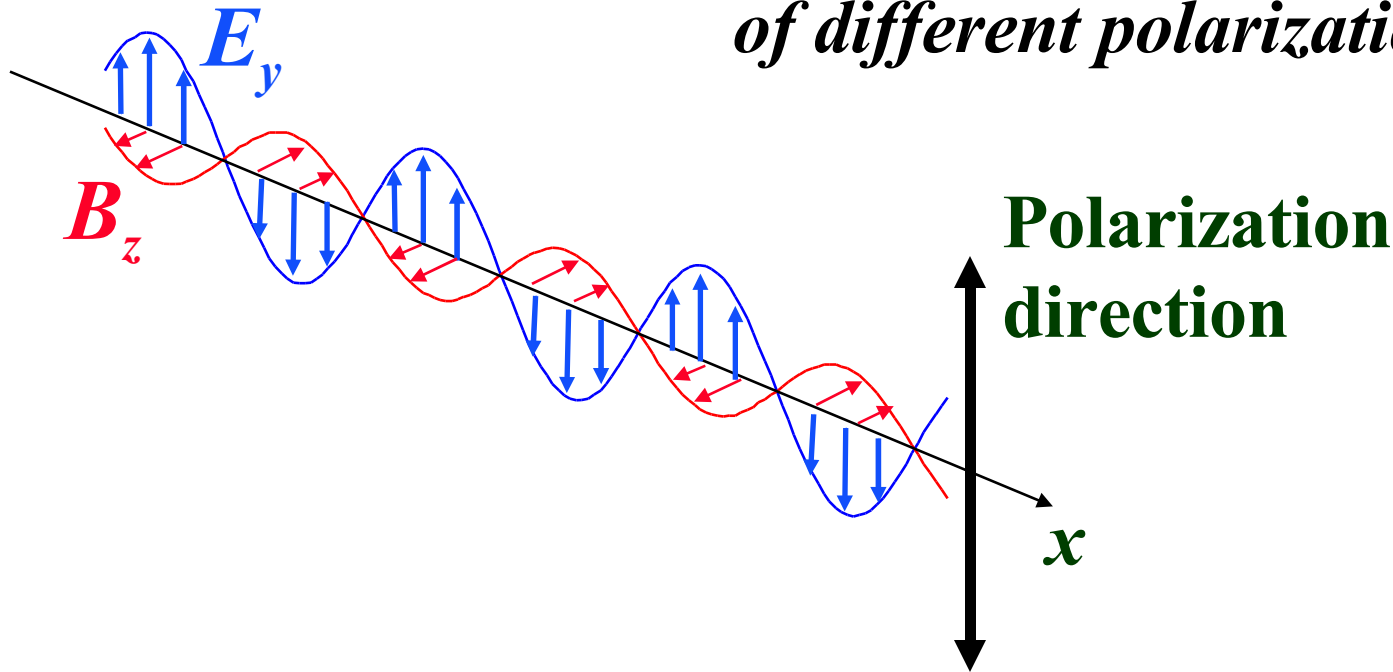
$$P = 2S/c = 2(1400 \text{ W/m}^2) / 3 \times 10^8 \text{ m/s} = 9.4 \times 10^{-6} \text{ N/m}^2$$

Hence $A=1000\text{N} / 9.4 \times 10^{-6} \text{ N/m}^2 = 1.0 \times 10^8 \text{ m}^2 = 100 \text{ km}^2$

Polarization

n

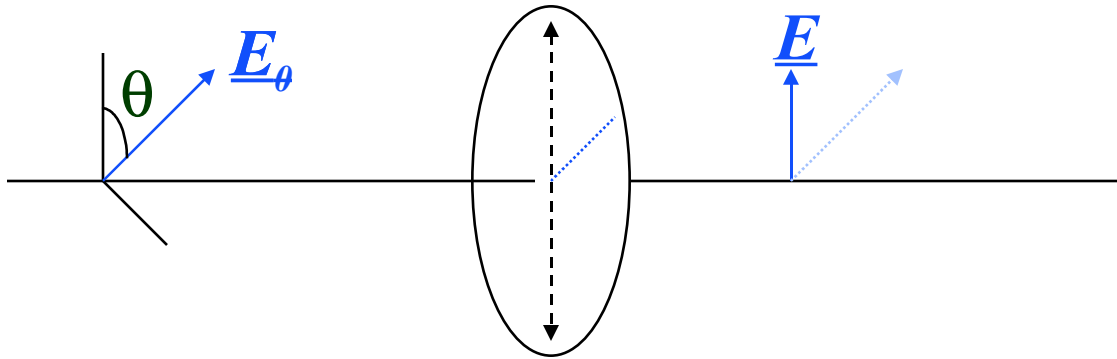
The direction of polarization of a wave is the direction of the electric field. Most light is randomly polarized, which means it contains a mixture of waves of different polarizations.



Polarization

n

A polarizer lets through light of only one polarization:



Transmitted light has its E in the direction of the polarizer's transmission axis.

$$E = E_0 \cos \theta$$

hence,

$$S = S_0 \cos^2 \theta$$

- Malus's Law