

Chapter 32

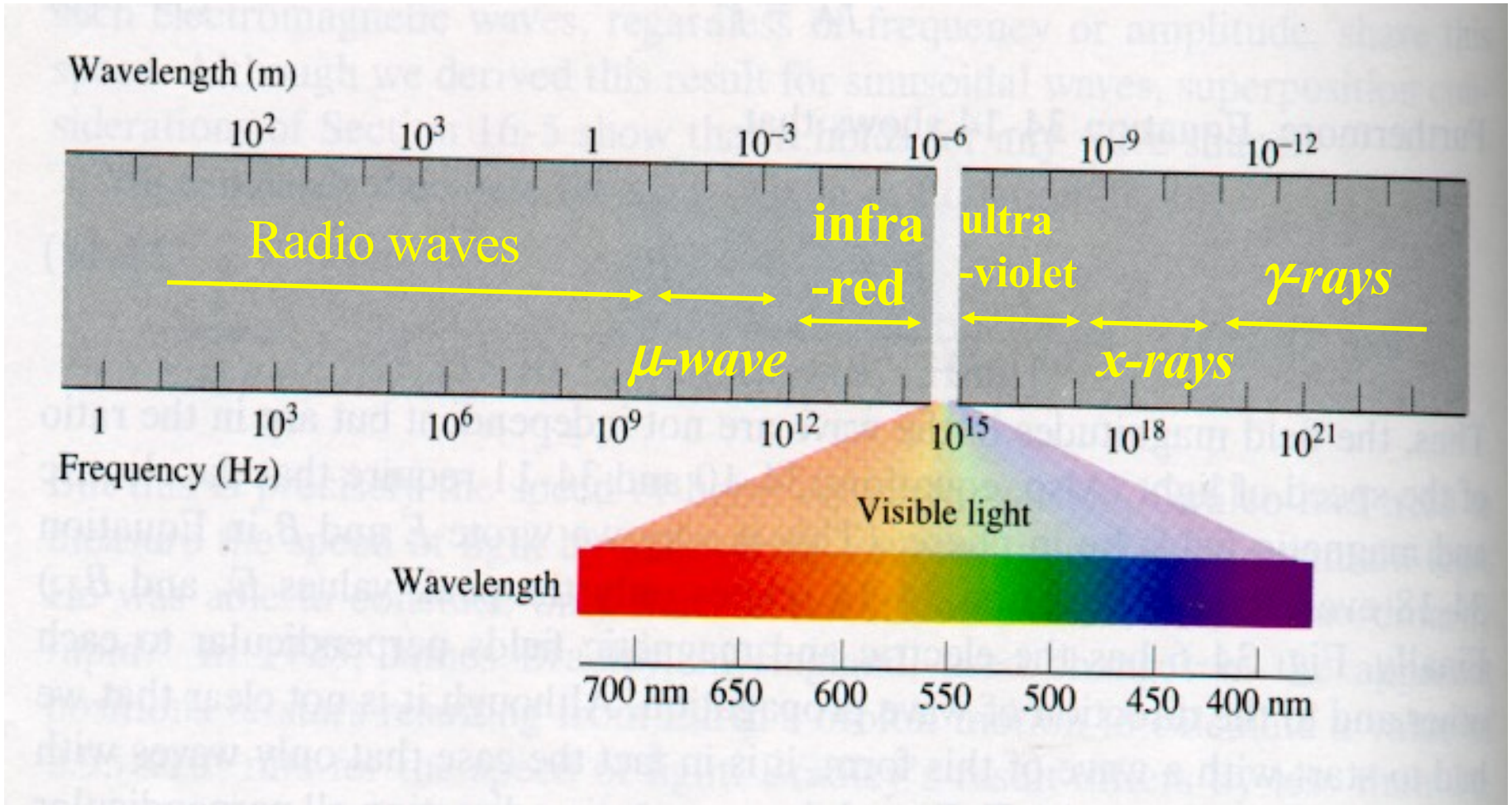
The Laws of Electromagnetism

Maxwell's Equations

Displacement Current

Electromagnetic Radiation

The Electromagnetic Spectrum



The Equations of Electromagnetism (at this point ...)

Gauss' Law for Electrostatics $\oint \underline{E} \cdot \underline{dA} = \frac{q}{\epsilon_0}$

Gauss' Law for Magnetism $\oint \underline{B} \cdot \underline{dA} = 0$

Faraday's Law of Induction $\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$

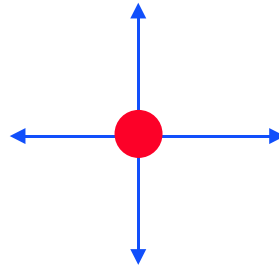
Ampere's Law $\oint \underline{B} \cdot \underline{dl} = \mu_0 I$

The Equations of Electromagnetism

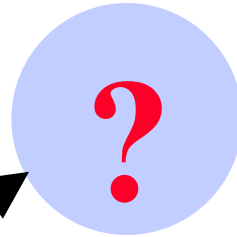
Gauss's Laws

..monopole..

1 $\oint \underline{E} \cdot \underline{dA} = \frac{q}{\epsilon_0}$



2 $\oint \underline{B} \cdot \underline{dA} = 0$



...there's no magnetic monopole....!!

The Equations of Electromagnetism

Faraday's Law

$$3 \quad \oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$

.. if you change a magnetic field you induce an electric field.....

Ampere's Law

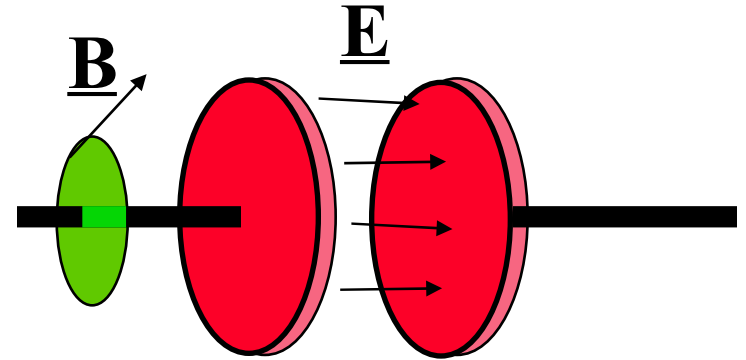
$$4 \quad \oint \underline{B} \cdot \underline{dl} = \mu_0 I$$

.....is the reverse true..?

...lets take a look at charge flowing into a capacitor...

...when we derived Ampere's Law we assumed constant current...

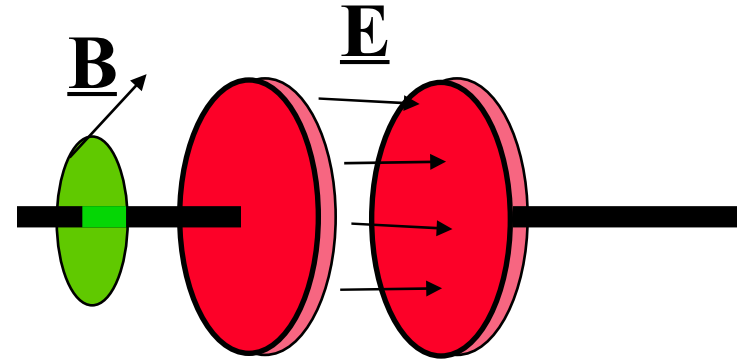
$$\oint \underline{B} \cdot \underline{dl} = \mu_0 I$$



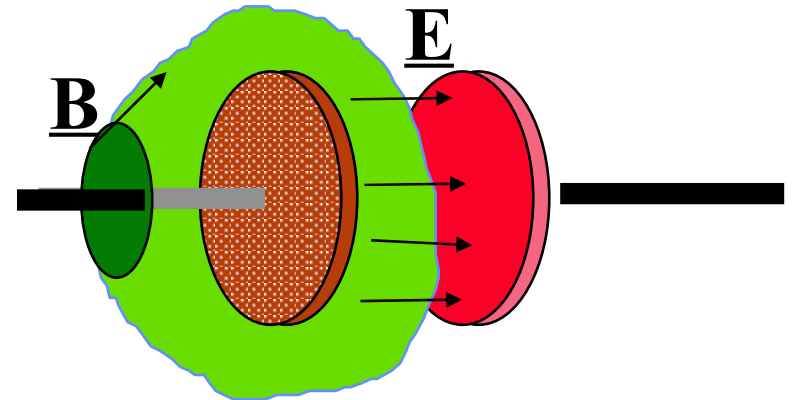
...lets take a look at charge flowing into a capacitor...

...when we derived Ampere's Law we assumed constant current...

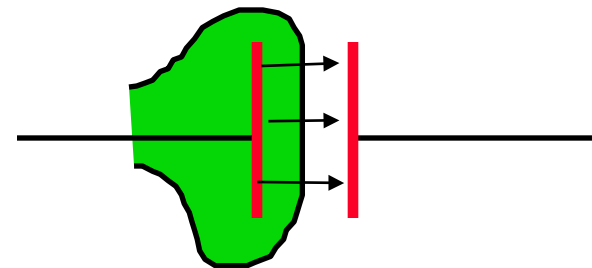
$$\oint \underline{B} \cdot d\underline{l} = \mu_0 I$$



.. if the loop encloses one plate of the capacitor..there is a problem ... $I = 0$

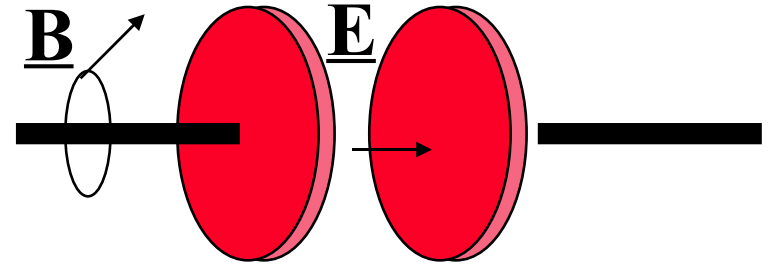


Side view: (Surface is now like a bag:)



*Maxwell solved this problem
by realizing that...*

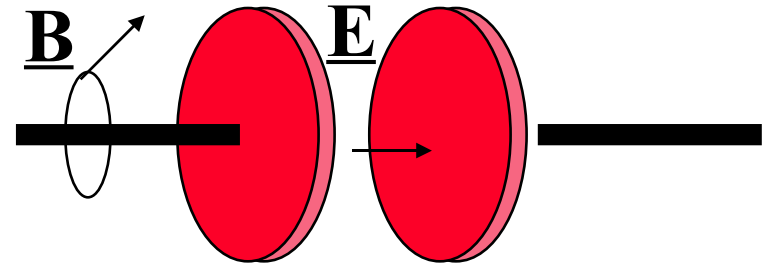
Inside the capacitor there must
be an induced magnetic field...



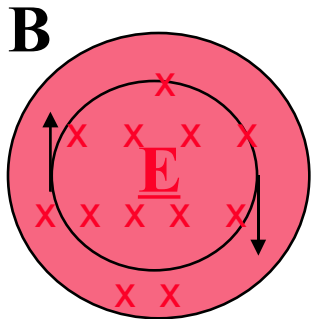
How?.

Maxwell solved this problem by realizing that...

Inside the capacitor there must be an induced magnetic field...



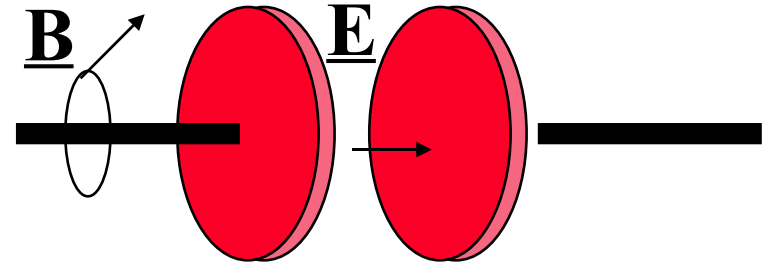
How?. Inside the capacitor there is a changing $\underline{E} \Rightarrow$



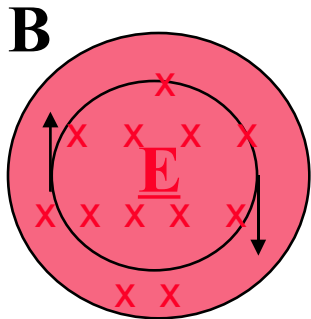
A changing
electric field
induces a
magnetic field

Maxwell solved this problem by realizing that...

Inside the capacitor there must be an induced magnetic field...



How?. Inside the capacitor there is a changing $\underline{E} \Rightarrow$



A changing
electric field
induces a
magnetic field

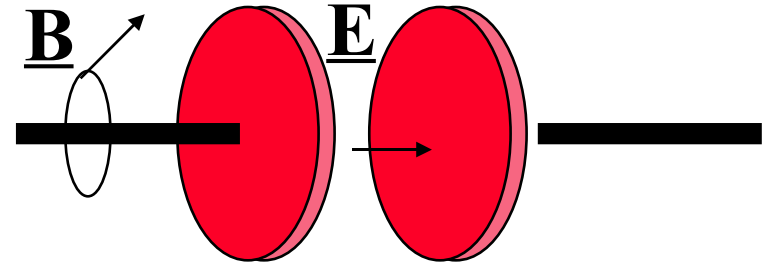
$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt} = \mu_0 I_d$$

where I_d is called the
displacement current

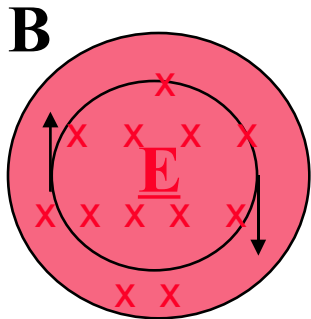
-

Maxwell solved this problem by realizing that...

Inside the capacitor there must be an induced magnetic field...



How?. Inside the capacitor there is a changing $E \Rightarrow$



A changing electric field induces a magnetic field

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt} = \mu_0 I_d$$

where I_d is called the displacement current

Therefore, Maxwell's revision of Ampere's Law becomes....

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Derivation of Displacement Current

For a capacitor, $q = \epsilon_0 EA$ and $I = \frac{dq}{dt} = \epsilon_0 \frac{d(EA)}{dt}$.

Now, the electric flux is given by EA , so: $I = \epsilon_0 \frac{d(\Phi_E)}{dt}$,
where this current , not being associated with charges, is
called the “Displacement current”, I_d .

Hence:
$$I_d = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Derivation of Displacement Current

For a capacitor, $q = \epsilon_0 EA$ and $I = \frac{dq}{dt} = \epsilon_0 \frac{d(EA)}{dt}$.

Now, the electric flux is given by EA , so: $I = \epsilon_0 \frac{d(\Phi_E)}{dt}$, where this current, not being associated with charges, is called the “Displacement Current”, I_d .

Hence:
$$I_d = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

and:
$$\oint \underline{B} \cdot \underline{dl} = \mu_0 (I + I_d)$$

$$\Rightarrow \oint \underline{B} \cdot \underline{dl} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Maxwell's Equations of Electromagnetism

Gauss' Law for Electrostatics $\oint \underline{E} \cdot \underline{dA} = \frac{q}{\epsilon_0}$

Gauss' Law for Magnetism $\oint \underline{B} \cdot \underline{dA} = 0$

Faraday's Law of Induction $\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$

Ampere's Law $\oint \underline{B} \cdot \underline{dl} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$

Maxwell's Equations of Electromagnetism in Vacuum (no charges, no masses)

Consider these equations in a vacuum.....

.....no mass, no charges. no currents.....

$$\oint \underline{E} \cdot \underline{dA} = \frac{q}{\epsilon_0} \longrightarrow \oint \underline{E} \cdot \underline{dA} = 0$$

$$\oint \underline{B} \cdot \underline{dA} = 0 \longrightarrow \oint \underline{B} \cdot \underline{dA} = 0$$

$$\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt} \longrightarrow \oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt} \longrightarrow \oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Maxwell's Equations of Electromagnetism in Vacuum (no charges, no masses)

$$\oint \underline{E} \cdot \underline{dA} = 0$$

$$\oint \underline{B} \cdot \underline{dA} = 0$$

$$\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Electromagnetic Waves

Faraday's law: $d\mathbf{B}/dt \longrightarrow$ electric field

Maxwell's modification of Ampere's law

$d\mathbf{E}/dt \longrightarrow$ magnetic field

$$\oint \underline{\mathbf{B}} \cdot \underline{d\mathbf{l}} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt} \qquad \oint \underline{\mathbf{E}} \cdot \underline{d\mathbf{l}} = - \frac{d\Phi_B}{dt}$$

These two equations can be solved simultaneously.

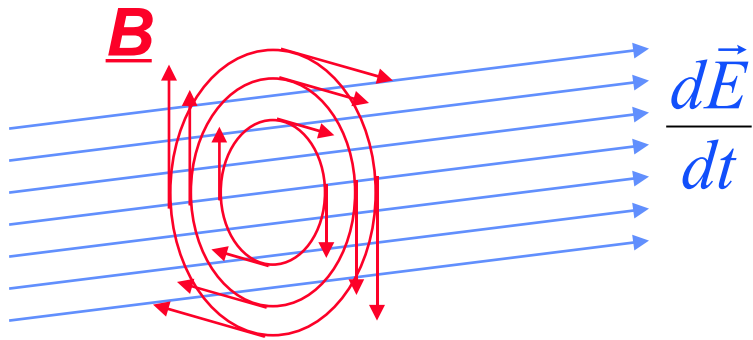
The result is:

$$\underline{\mathbf{E}}(\mathbf{x}, t) = E_p \sin(\mathbf{kx} - \omega t) \hat{\mathbf{j}}$$

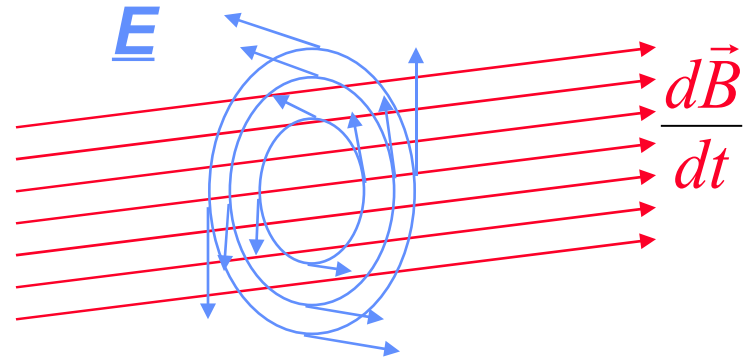
$$\underline{\mathbf{B}}(\mathbf{x}, t) = B_p \sin(\mathbf{kx} - \omega t) \hat{\mathbf{z}}$$

Electromagnetic Waves

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$



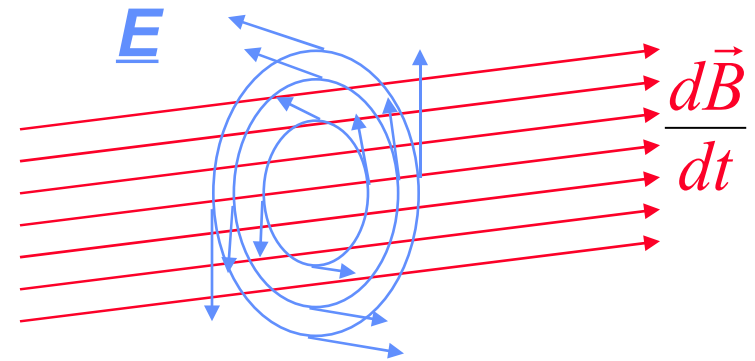
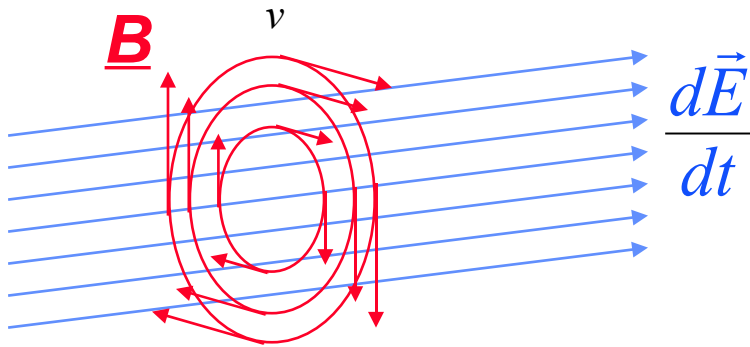
$$\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$



Electromagnetic Waves

$$\oint \underline{B} \cdot \underline{dl} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

$$\oint \underline{E} \cdot \underline{dl} = - \frac{d\Phi_B}{dt}$$



Special case..PLANE WAVES...

$$\vec{E} = E_y(x,t)\hat{j} \quad \vec{B} = B_z(x,t)\hat{k}$$

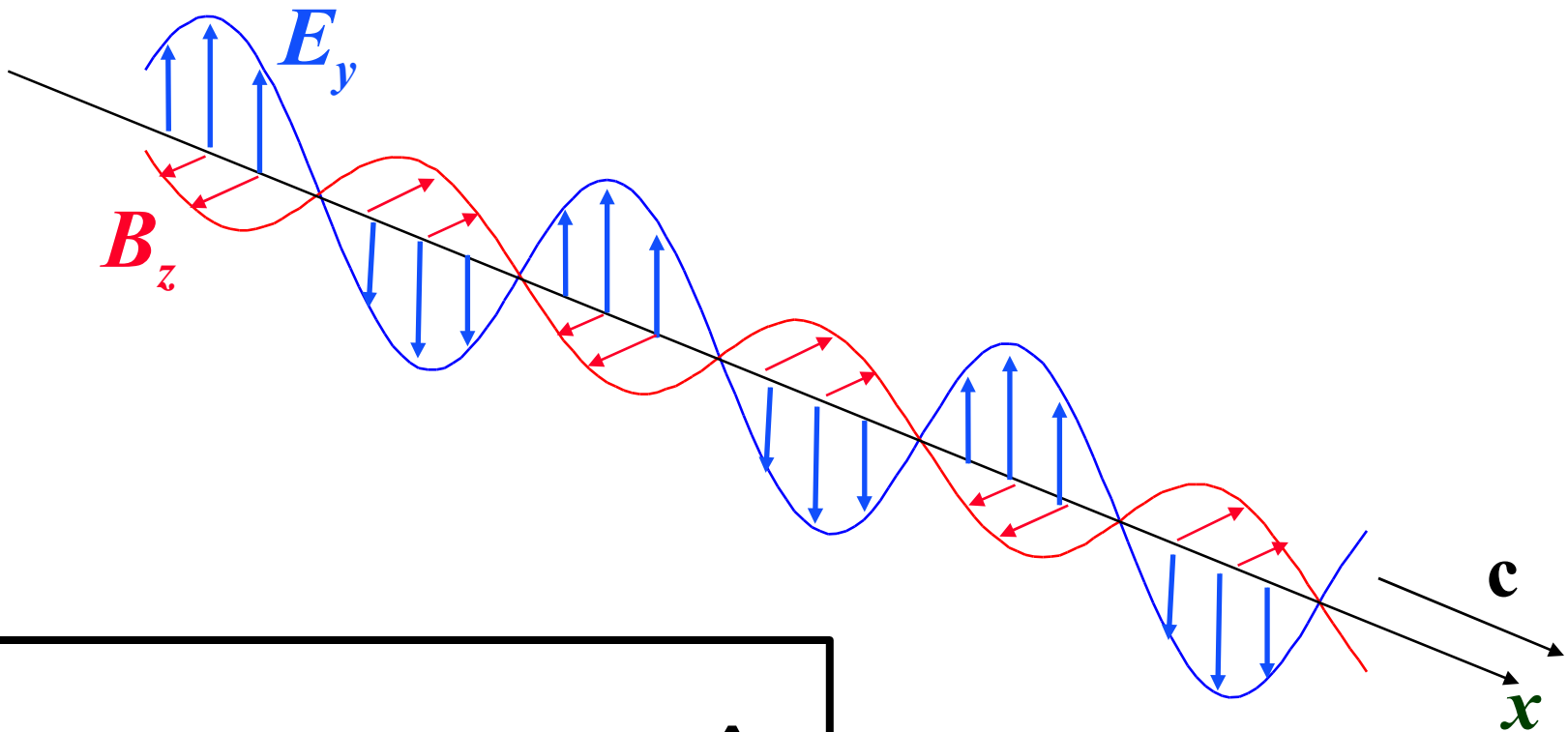
satisfy the wave equation

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2}$$

Maxwell's Solution

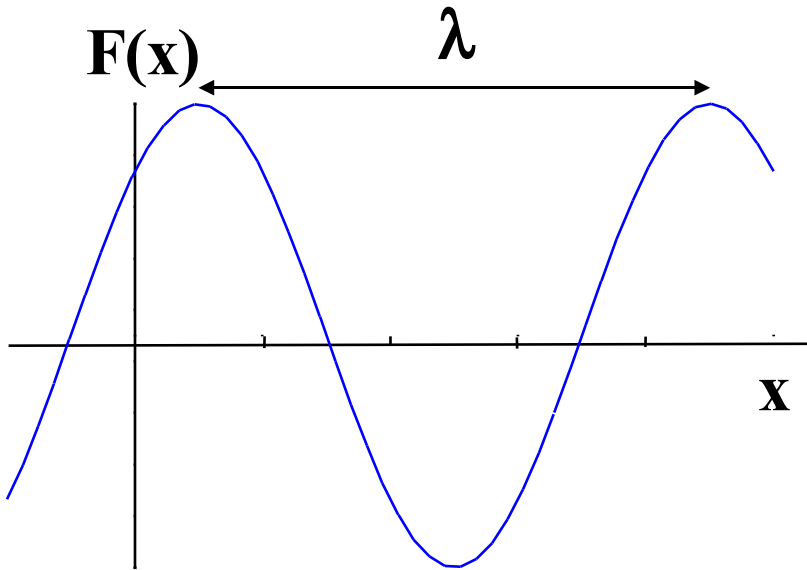
$$\psi = A \sin(\omega t + \phi)$$

Plane Electromagnetic Waves



$$\underline{\mathbf{E}}(\mathbf{x}, t) = E_p \sin (kx - \omega t) \hat{\mathbf{j}}$$

$$\underline{\mathbf{B}}(\mathbf{x}, t) = B_p \sin (kx - \omega t) \hat{\mathbf{z}}$$



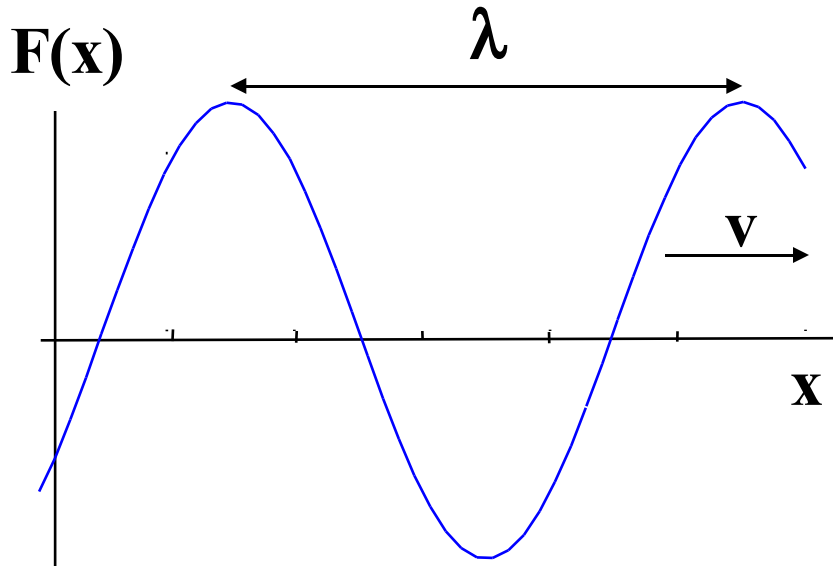
Static wave

$$F(x) = F_p \sin (kx + \phi)$$

$$k = 2\pi / \lambda$$

k = wavenumber

λ = wavelength



Moving wave

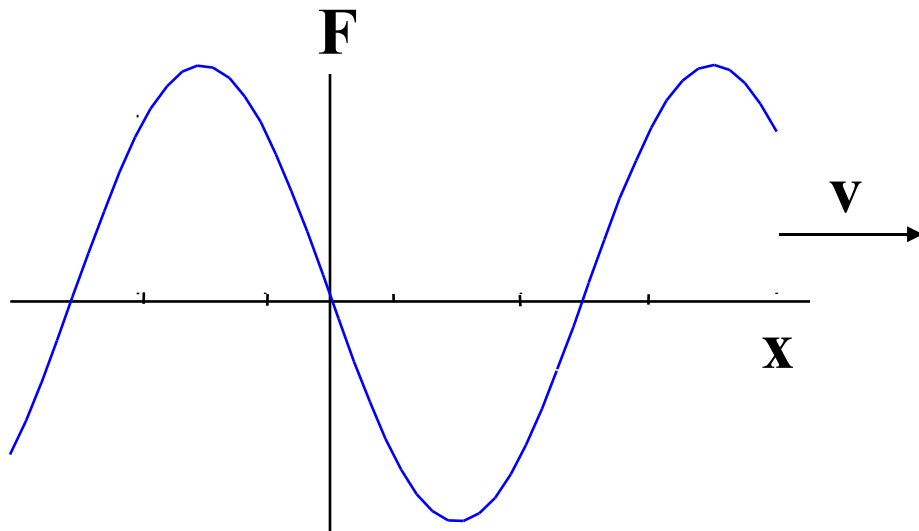
$$F(x, t) = F_p \sin (kx - \omega t)$$

$$\omega = 2\pi / f$$

ω = angular frequency

f = frequency

$$v = \omega / k$$



Moving wave

$$F(x, t) = F_p \sin (kx - \omega t)$$

What happens at $x = 0$ as a function of time?

$$F(0, t) = F_p \sin (-\omega t)$$

$$\text{For } x = 0 \text{ and } t = 0 \Rightarrow F(0, 0) = F_p \sin (0)$$

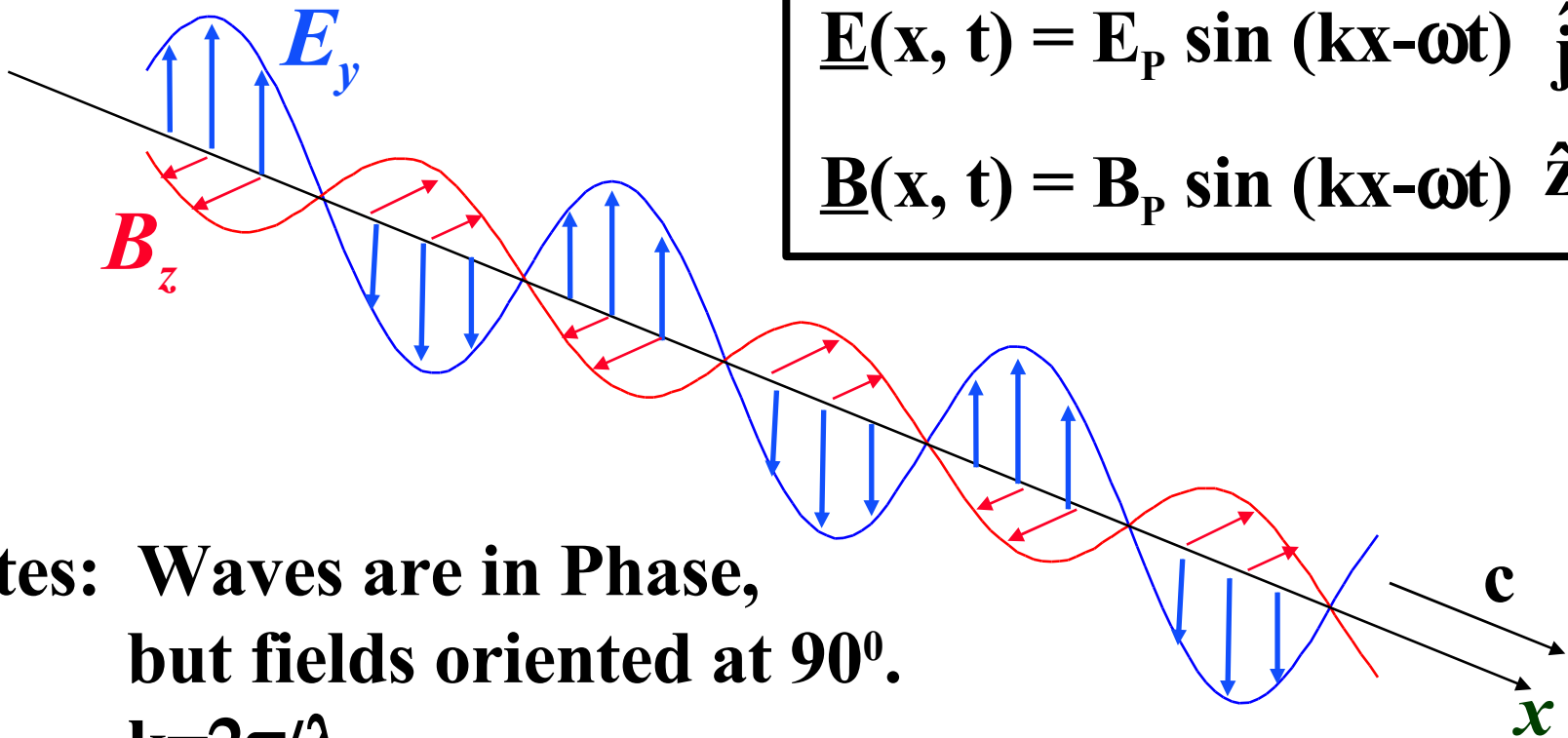
$$\text{For } x = 0 \text{ and } t = t \Rightarrow F (0, t) = F_p \sin (0 - \omega t) = F_p \sin (-\omega t)$$

This is equivalent to: $kx = -\omega t \Rightarrow x = -(\omega/k) t$

$F(x=0)$ at time t is the same as $F[x=-(\omega/k)t]$ at time 0

The wave moves to the right with speed ω/k

Plane Electromagnetic Waves



**Notes: Waves are in Phase,
but fields oriented at 90° .**

$$k = 2\pi/\lambda.$$

Speed of wave is $c = \omega/k$ ($= f\lambda$)

$$c = 1 / \sqrt{\epsilon_0 \mu_0} = 3 \times 10^8 \text{ m/s}$$

At all times $E = cB$.